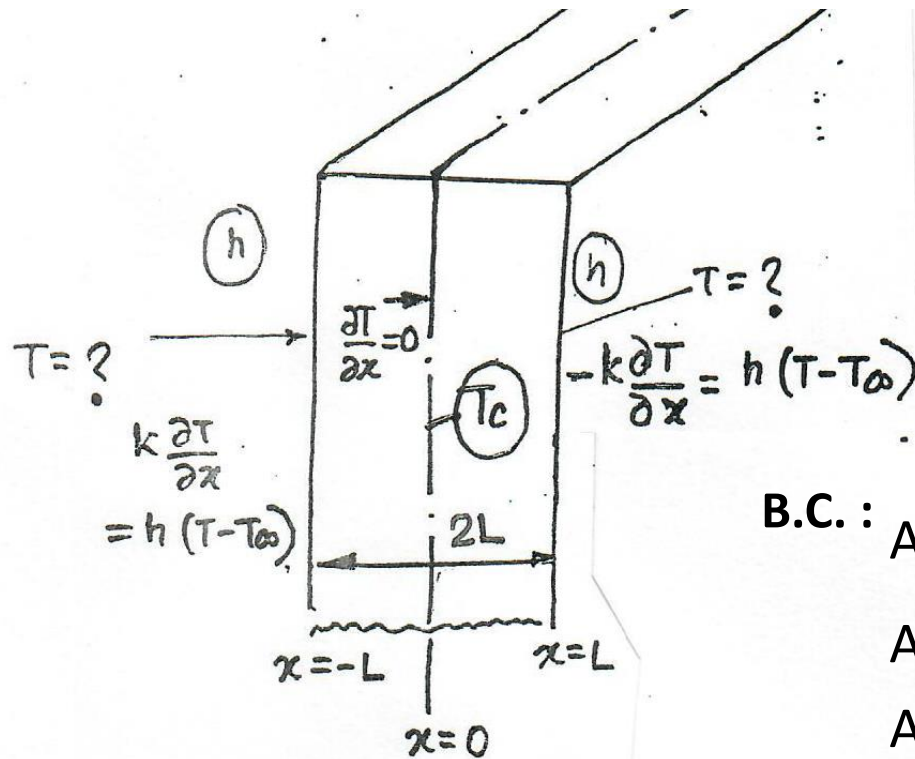


Heat and Mass transfer

UNSTEADY STATE CONDUCTION IN **AN INFINITE SLAB**

TRANSIENT CONDUCTION IN RECTANGULAR GEOMETRY WITH CONVECTIVE BOUNDARY CONDITION



- Local ambience is at $T = T_\infty$.
- Temperature at the surface is not known.
- h is known.

G.E.:
$$\frac{\partial T}{\partial t} = \alpha \frac{\partial^2 T}{\partial x^2}$$

B.C.: At $x = 0$, $\frac{\partial T}{\partial x} = 0$, for $t > 0$

At $x = L$, $-k \frac{\partial T}{\partial x} = h(T - T_\infty)$, for $t > 0$

At $t = 0$, $T = T_o$, for $0 \leq x \leq L$

Put, $\theta = (T - T_\infty)$,

G.E.:
$$\frac{\partial \theta}{\partial t} = \alpha \frac{\partial^2 \theta}{\partial x^2}$$

B.C.: At $x = 0$, $\frac{\partial \theta}{\partial x} = 0$, for $t > 0$

At $x = L$, $-k \frac{\partial \theta}{\partial x} = h\theta$, for $t > 0$

At $t = 0$, $\theta = \theta_o$, for $0 \leq x \leq L$

Let, $\theta = X(x)\Gamma(t)$, and the governing equation will be

$$\alpha X''\Gamma = X \dot{\Gamma}$$

$$\text{or, } \frac{\dot{\Gamma}}{\alpha\Gamma} = \frac{X''}{X} = \text{constant} = \pm\lambda^2$$

As it was shown in the earlier problem, the constant should be $-\lambda^2$ and the solution for Γ , is given by

$$\Gamma = c_1 e^{-\alpha\lambda^2 t}$$

Where, c_1 is any arbitrary constant.

Now R.H.S. can be written as: $\frac{X''}{X} = -\lambda^2$

Solution is, $X = c_2 \cos\lambda x + c_3 \sin\lambda x$, where c_2 and c_3 are arbitrary constants.

$$\text{Now, } \theta = X(x)\Gamma(t) = e^{-\alpha\lambda^2 t}[A \cos\lambda x + B \sin\lambda x] \quad (1)$$

$$\frac{\partial\theta}{\partial x} = e^{-\alpha\lambda^2 t}[-A\lambda \sin\lambda x + B\lambda \cos\lambda x]$$

At $x = 0$, $\frac{\partial \theta}{\partial x} = 0$ } Substituting this boundary condition in the above equation
for $t > 0$ yields, $B = 0$

Therefore, $\theta = e^{-\alpha \lambda^2 t} A \cos \lambda x$

$$-k \frac{\partial \theta}{\partial x} = +k \lambda A \sin \lambda x e^{-\alpha \lambda^2 t}$$

$$-k \frac{\partial \theta}{\partial x} \Big|_{x=L} = h \theta$$

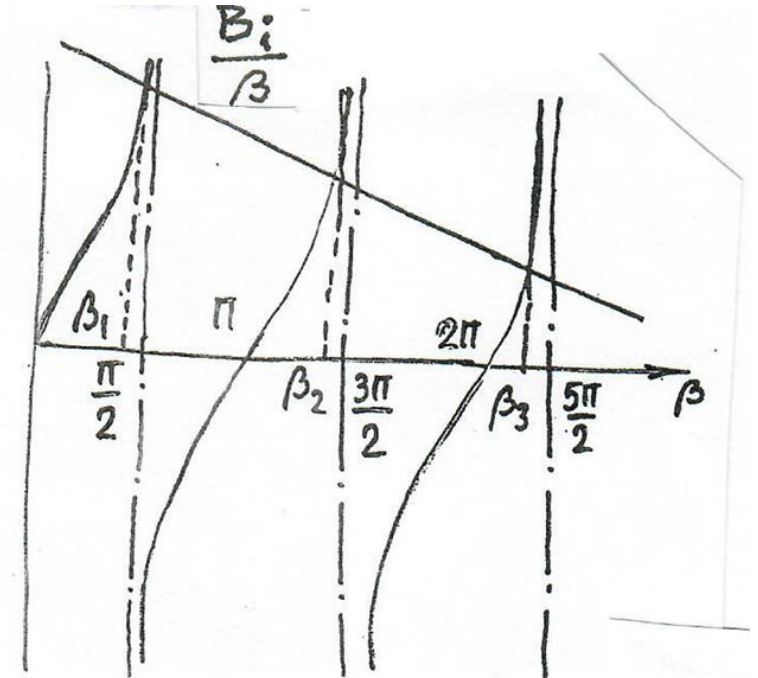
$$k \lambda A \sin \lambda L e^{-\alpha \lambda^2 t} = h e^{-\alpha \lambda^2 t} A \cos \lambda L$$

$$\text{or, } \lambda \tan \lambda L = \frac{h}{k}$$

$$\text{or, } \lambda L \tan \lambda L = \frac{hL}{k} = Bi \quad (2)$$

where, **Bi** is **Biot Number**.

$$\text{Or, } \beta \tan \beta = Bi, \text{ where } \beta = \lambda L \quad (3)$$



This equation is a transcendental equation and it has an infinite number of roots, the eigen values β , viz, $\beta_1, \beta_2, \beta_3, \dots, \beta_n$ are shown graphically (points of intersection). These are also available in a tabular form.

Once $\beta_n = \lambda_n L$, is known λ_n is known, and we get the solution as,

$$\theta = \sum_{n=1}^{\infty} A_n \cos \lambda_n x e^{-\alpha \lambda_n^2 t}$$

The initial condition is : at $t = 0$, $\theta = \theta_o$

$$\theta_o = \sum_{n=1}^{\infty} A_n \cos \lambda_n x$$

$$\int_0^L \theta_o \cos \lambda_m x \, dx = \int_0^L \{A_n \cos \lambda_n x\} \cos \lambda_m x \, dx$$

For $m = n$,

$$\int_0^L \theta_o \cos \lambda_n x \, dx = \int_0^L \Sigma A_n \cos^2 \lambda_n x \, dx$$

or,

$$\left[\frac{\theta_o \sin \lambda_n x}{\lambda_n} \right]_0^L = \frac{A_n}{2} \int_0^L [1 + \cos 2\lambda_n x] dx$$

For $m = n$,

$$\int_0^L \theta_o \cos \lambda_n x \, dx = \int_0^L \Sigma A_n \cos^2 \lambda_n x \, dx$$

or,

$$\left[\frac{\theta_o \sin \lambda_n x}{\lambda_n} \right]_0^L = \frac{A_n}{2} \int_0^L [1 + \cos 2\lambda_n x] dx$$

$$\left[\frac{\theta_o \sin \lambda_n L}{\lambda_n} \right] = \frac{A_n}{2} \left[x + \frac{\sin 2\lambda_n x}{2\lambda_n} \right]_0^L$$

$$A_n = \frac{4 \theta_o \sin \lambda_n L}{[2 \lambda_n L + \sin 2 \lambda_n L]}$$

$$A_n = \frac{4 \theta_o \sin \beta_n}{2 \beta_n + \sin 2 \beta_n} \quad (4)$$

Where, $\beta_n = \lambda_n L$

Now,

$$\theta = \sum_{n=1}^{\infty} \frac{4\theta_o \sin \beta_n}{2\beta_n + \sin 2\beta_n} \cos \left(\frac{\beta_n x}{L} \right) e^{-\alpha \left(\frac{\beta_n}{L} \right)^2 t} \quad (5)$$

We know, $\beta_n = \lambda_n L$

We also observe that, $\frac{\theta}{\theta_o} = f \left(\lambda_n L, \frac{\alpha t}{L^2}, \frac{x}{L} \right)$

Where, $\frac{\alpha t}{L^2} = \mathbf{Fo}$, (**Fourier Number**)

From **(2)**, we have also seen that,

$$\lambda L = f \left(\frac{hL}{k} \right) = f(Bi)$$

We may say that the dimensionless temperature distribution is a function of Bi , Fo and $\frac{x}{L}$

Heisler showed that the infinite series in equation (5) converges rapidly and that for values of $\frac{\alpha t}{L^2}$ greater than 0.2, all terms after the first can be neglected.

Thus one needs to know only the first eigen value in the equation(5).Using this fact, equation(5) reduces to

$$\frac{\theta}{\theta_o} = \frac{4 \sin \lambda_1 L}{2\lambda_1 L + \sin^2 (2\lambda_1 L)} \cos \left(\lambda_1 L \cdot \frac{x}{L} \right) e^{-(\lambda_1 L)^2 \frac{\alpha t}{L^2}} \quad (6)$$

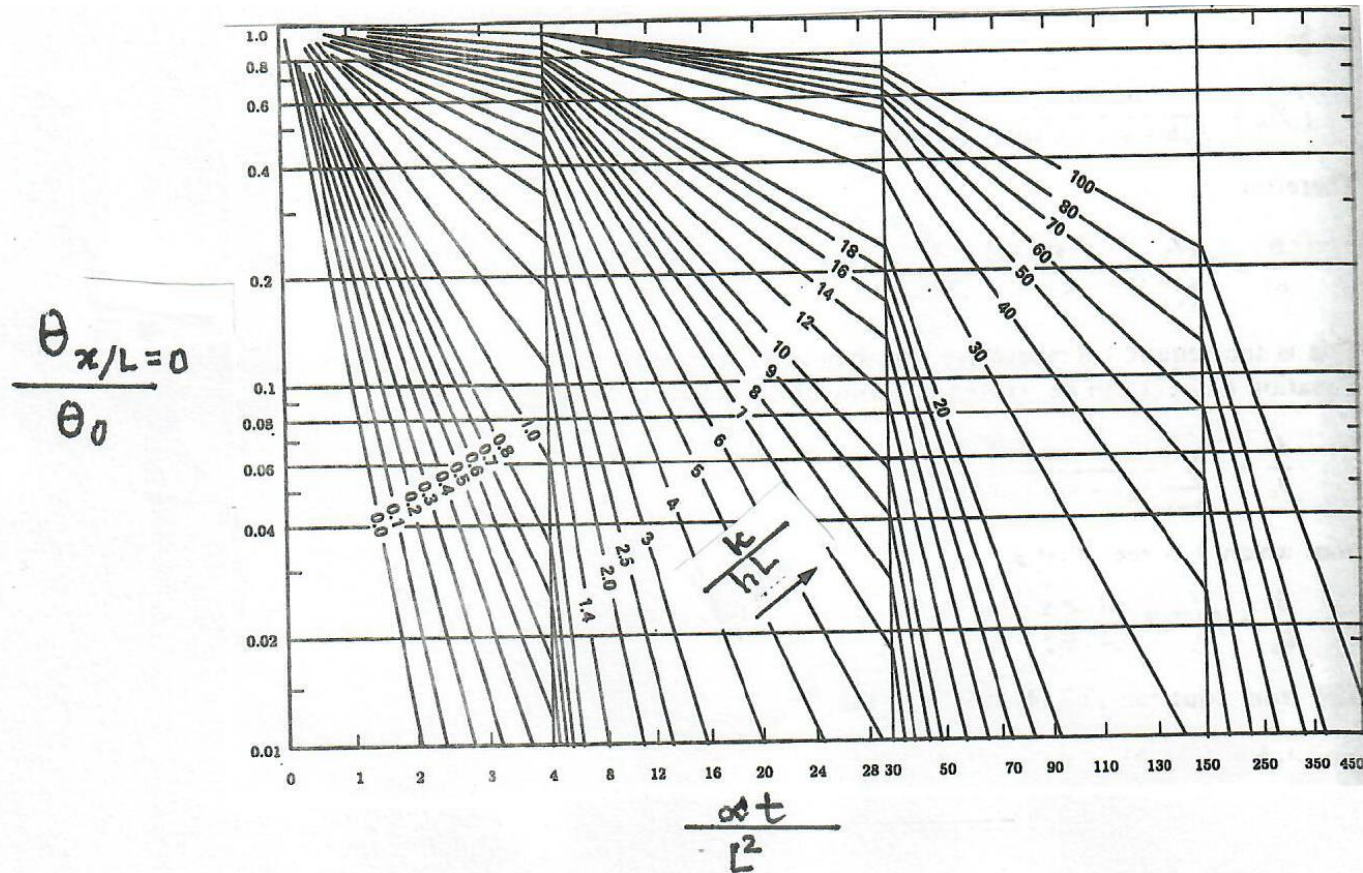
Setting, $\frac{x}{L} = 0$, we have the dimensionless temperature at the centre-line given by

$$\frac{\theta_{x/L=0}}{\theta_o} = \frac{4 \sin \lambda_1 L}{2\lambda_1 L + \sin^2 (2\lambda_1 L)} e^{-(\lambda_1 L)^2 \frac{\alpha t}{L^2}} \quad (7)$$

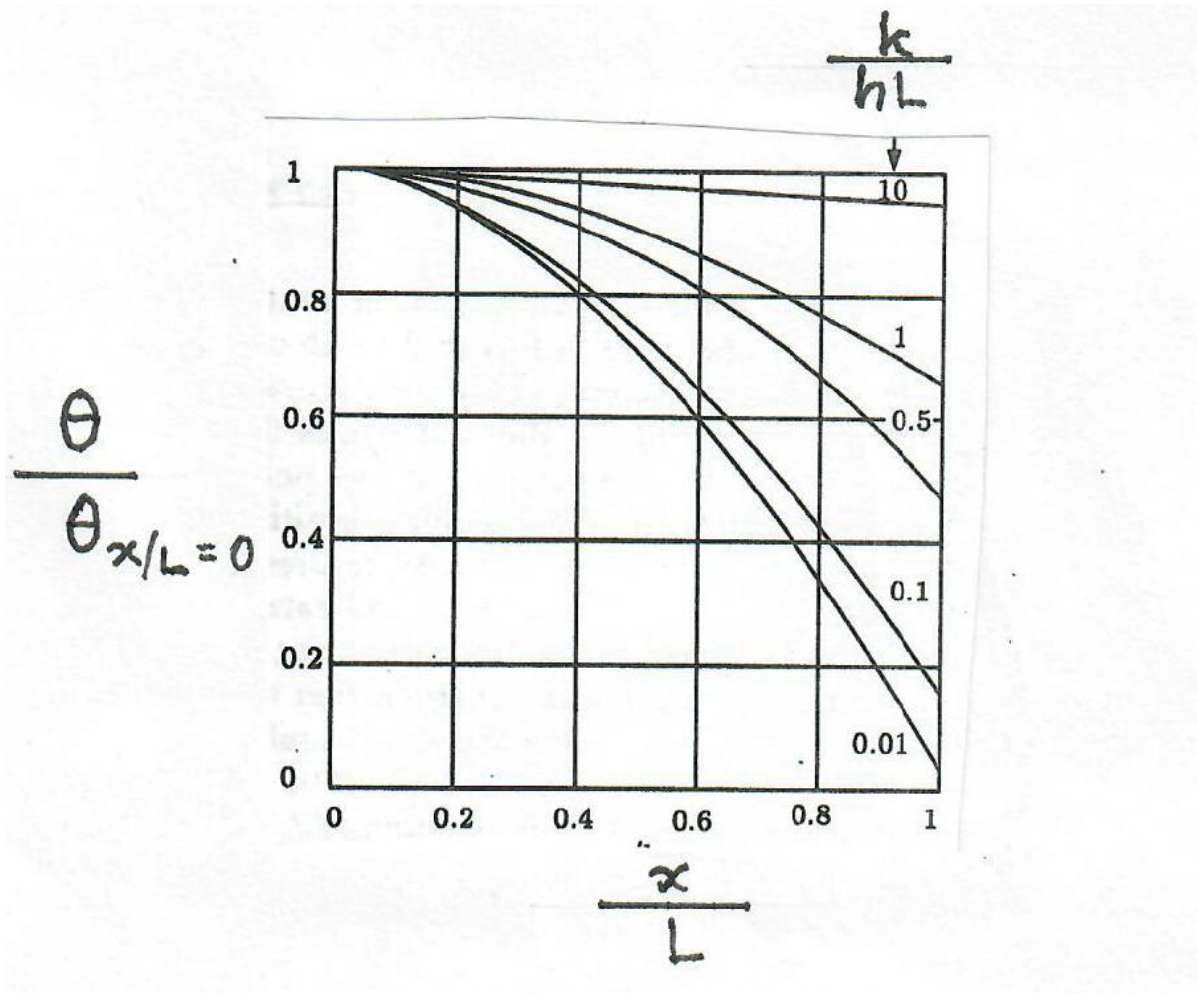
Dividing equation (6) by equation(7), we have

$$\frac{\theta}{\theta_{x/L=0}} = \cos \left(\lambda_1 L \cdot \frac{x}{L} \right) \quad (8)$$

- Thus the ratio of the temperature at any point in the slab to the temperature at the centre-line is independent of time.
- Consequently the complete temperature history in the slab can be drawn in two charts, called **Heisler Charts** (named after M.P. Heisler).
- In one chart, the dimensionless centre-line temperature ($\theta_{x/L=0}$) θ_o is plotted against the Fourier number Fo , with the inverse of the **Biot Number** Bi as the parameter.



- In the second chart, the dimensionless ratio of the temperature at any point in the slab and the centre-line temperature is plotted against the dimensionless position.



The corresponding analytical solution for the one- dimensional unsteady state temperature distributions in a long solid cylinder or a solid sphere, initially at a uniform temperature and subsequently subjected to a different ambient temperatures, can be obtained in a similar fashion. Heisler charts are also available to determine the temperature distribution graphically.

Thank you